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## Numerical methods for hyperbolic systems

### Part 1 of correction exercise sheet 2: Galerkin discontinuous for advection equation

**Exercise 1** We consider the Lax-Wendroff scheme

$$\frac{u_j^{n+1} - u_j^n}{\Delta t} + \frac{a}{2\Delta x}(u_{j+1}^n - u_{j-1}^n) - \frac{a^2\Delta t}{2\Delta x^2}(u_{j+1}^n - 2u_j^n + u_{j-1}^n) = 0 \quad (1)$$

with  $\Delta t$  the time step,  $\Delta x$  the step mesh and  $u_j^n$  the approximation to  $u(n\Delta t, j\Delta x)$  where  $n \in \mathbb{N}$ ,  $j \in \mathbb{N}$ .

1. Study the  $L^2$  stability.

## $L^2$ stability with continuous Fourier transform

The Neumann analysis is based on Fourier transformation. We introduce the Fourier coefficient  $\hat{u}^n(k)$ . The scheme written in the Fourier space is

$$\hat{u}^{n+1}(k) = \hat{u}^n(k) - \frac{a\Delta t}{2\Delta x} \left( e^{i\zeta} \hat{u}^n(k) - e^{-i\zeta} \hat{u}^n(k) \right) + \frac{a^2\Delta t}{2\Delta x^2} \left( 2\hat{u}^n(k) - e^{i\zeta} \hat{u}^n(k) - e^{-i\zeta} \hat{u}^n(k) \right)$$

with  $\zeta = 2\pi k\Delta x$ . After simplification we obtain

$$\hat{u}^{n+1}(k) = \left( 1 - \left( \frac{\lambda}{2}(e^{i\zeta} - e^{-i\zeta}) + \frac{\lambda^2}{2}(2 - e^{i\zeta} - e^{-i\zeta}) \right) \right) \hat{u}^n(k),$$

with  $\lambda = \frac{a\Delta t}{\Delta x}$ . Now we use  $e^{\pm i\zeta} = \cos(\zeta) \pm i \sin(\zeta)$ . We obtain

$$\hat{u}^{n+1}(k) = A(k)\hat{u}^n(k),$$

with  $A(k) = 1 - (\lambda i \sin(\zeta) + \lambda^2(1 - \cos(\zeta)))$ .

To prove the stability we must prove  $|\hat{u}^{n+1}(k)|^2 = |A(k)|^2|\hat{u}^n(k)|^2$  with  $|A(k)| \leq 1$ . By definition of the complex module we have

$$|A(k)|^2 = (1 - \lambda^2(1 - \cos(\zeta)))^2 + \lambda^2 \sin^2(\zeta).$$

Expanding the previous expression we obtain

$$|A(k)|^2 = 1 - 2\lambda^2(1 - \cos(\zeta)) + \lambda^4(1 - \cos(\zeta))^2 + \lambda^2(1 - \cos^2(\zeta)).$$

Now we simplify to obtain

$$|A(k)|^2 = 1 - \lambda^2(1 - 2\cos(\zeta) + 2\cos^2(\zeta)) + \lambda^4(1 - \cos(\zeta))^2.$$

$$|A(k)|^2 = 1 - \lambda^2(1 - \cos(\zeta))^2 + \lambda^4(1 - \cos(\zeta))^2 = 1 + (\lambda^4 - \lambda^2)(1 - \cos(\zeta))^2.$$

If  $\lambda \leq 1$  the term  $(\lambda^4 - \lambda^2) \leq 0$  and  $|A(k)| \leq 1$ . Consequently the scheme is stable

$$\frac{a\Delta t}{\Delta x} \leq 1.$$

## $L^2$ stability with discrete Fourier transform

The Lax wendroff scheme is defined by

$$\mathbf{U}^{n+1} = C\mathbf{U}^n,$$

with  $C$  a circulant matrix defined by the coefficients  $c_0 = 1 - \alpha^2$ ,  $c_{n-1} = \frac{1}{2}(\alpha^2 + \alpha)$ ,  $c_1 = \frac{1}{2}(\alpha^2 - \alpha)$  and  $\alpha = \frac{a\Delta t}{\Delta x}$ . Apply the Fourier transform we obtain

$$P^*\mathbf{U}^{n+1} = P^*C P P^*\mathbf{U}^n,$$

$$\hat{\mathbf{U}}^{n+1} = \Lambda_C \hat{\mathbf{U}}^n.$$

Consequently we obtain  $\|\hat{\mathbf{U}}^{n+1}\| \leq \|\Lambda_C\| \|\hat{\mathbf{U}}^n\|$  and the scheme is stable if  $\max_k |\lambda_C^k| \leq 1$ . In this case the eigenvalues of  $\Lambda_C$  are defined by

$$\lambda_C^k = 1 - \alpha^2 + \frac{1}{2}(\alpha^2 - \alpha)e^{\frac{2i\pi k}{n}} + \frac{1}{2}(\alpha^2 + \alpha)e^{\frac{2i\pi k(n-1)}{n}},$$

$$\lambda_C^k = 1 - \alpha^2 + \frac{1}{2}(\alpha^2 - \alpha)e^{\frac{2i\pi k}{n}} + \frac{1}{2}(\alpha^2 + \alpha)e^{\frac{-2i\pi k}{n}}.$$

We use the fact that  $e^{2i\pi k} = 1$ . Using  $e^{\frac{2i\pi k}{n}} = \cos(\frac{2\pi k}{n}) \pm i \sin(\frac{2\pi k}{n})$  we obtain  $\lambda_C^k = A(k)$ . To finish we use the end of the previous proof.

2. Prove that the consistency error associated to the Lax-Wendroff scheme is  $O(\Delta x^2 + \Delta t^2)$  (use the fact that  $\partial_{tt}u - a^2\partial_{xx}u = 0$ ).

Before we study the consistency error, we study the exact solution.  $u$  is solution of  $\partial_t u + a\partial_x u = 0$ . Taking the derivative of the equation we prove  $u$  is solution to  $\partial_{tt}u + a\partial_{t,x}u = 0$  and  $\partial_{t,x}u + a\partial_{xx}u = 0$ . Consequently  $u$  is solution of

$$\partial_{tt}u - a^2\partial_{xx}u = 0. \quad (2)$$

Now we define  $u(x_j, t_n)$  the exact solution. Using the second order Taylor expansion in time and first order expansion in space (as previously), we obtain the following consistency error

$$\begin{aligned} E &= \partial_t u(x_j, t_n) + \frac{\Delta t}{2}\partial_{tt}u(x_j, t_n) + O(\Delta t^2) + a\partial_x u(x_j, t_n) \\ &\quad - \Delta t \frac{a^2}{2}\partial_{xx}u(x_j, t_n) + O(\Delta x^2 + \Delta x^2\Delta t). \end{aligned}$$

We obtain

$$\begin{aligned} E &= \partial_t u(x_j, t_n) + \frac{\Delta t}{2}(\partial_{tt}u(x_j, t_n) - \Delta t \frac{a^2}{2}\partial_{xx}u(x_j, t_n)) + O(\Delta t^2) + a\partial_x u(x_j, t_n) \\ &\quad - \Delta t \frac{a^2}{2}\partial_{xx}u(x_j, t_n) + O(\Delta x^2 + \Delta x^2\Delta t) \end{aligned}$$

Since  $u(x_j, t_n)$  is a solution of  $\partial_t u + c\partial_x u = 0$  and (2) we obtain  $E = O(\Delta x^2 + \Delta t^2)$ .