

Several numerical models for the simulation of bubble oscillations

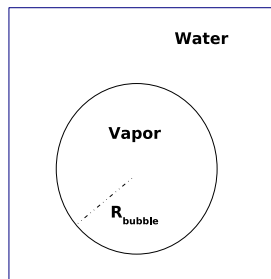
Hélène Mathis

IRMA, ULP, Strasbourg

25 janvier 2008

Joint work with M. Bachmann, P. Helluy, S. Müller

Introduction



- $p_{\text{bubble}} \ll p_{\text{liquid}}$
⇒ Oscillations and/or collapse
- Several models (energy model / isothermal model, bifluid / monofluid...)
- Different numerical approximations

Plan

- 1 Studied models
 - Euler system, stiffened gas law
 - Isothermal Euler system
 - Boundary conditions and initial data
- 2 Numerical scheme
- 3 Numerical results

Plan

1 Studied models

- Euler System
- Isothermal Euler equations
- Boundary conditions and initial data

2 Numerical scheme

3 Numerical results

Euler System

Unknowns:

- Density: $\rho(r, t)$
- Radial velocity: $v(r, t)$
- Internal energy: $e(r, t)$
- Fraction of gas: $\varphi(r, t)$

Pressure of the mixture: $p = p(\rho, e, \varphi)$

$$p(\rho, e, \varphi) = \begin{cases} (\gamma_1 - 1)\rho e & \text{if } \varphi = 1, \\ (\gamma_2 - 1)\rho e - \gamma_2 \pi_2 & \text{if } \varphi = 0. \end{cases} \quad (1)$$

with $\gamma_1 = 1.4$, $\gamma_2 = 1.1$ and $\pi_2 = \rho_0 c_0^2 / \gamma_2$.

Euler System

Mass conservation of gas:

$$(r^{d-1} \rho \varphi)_t + (r^{d-1} \rho \varphi u)_r = 0,$$

Mass conservation of liquid:

$$(r^{d-1} \rho (1 - \varphi))_t + (r^{d-1} \rho (1 - \varphi) u)_r = 0,$$

Momentum conservation:

$$(r^{d-1} \rho u)_t + (r^{d-1} (\rho u^2 + p))_x = (d - 1) p r^{(d-2)}, \quad (2)$$

Conservation of the total energy:

$$(r^{d-1} \rho E)_t + (r^{d-1} (\rho E + p) u)_x = 0,$$

Total energy: $E = e + u^2/2$.

with $d = 1, 2$ or 3

Isothermal Euler equations

φ is constant on the whole domain

Unknowns:

- Density: $\rho(r, t)$
- Radial velocity: $v(r, t)$

$$\begin{aligned}(r^{d-1}\rho)_t + (r^{d-1}\rho v)_r &= 0, \\ (r^{d-1}\rho v)_t + (r^{d-1}(\rho v^2 + p))_x &= (d-1)pr^{(d-2)}\end{aligned}\tag{3}$$

with $d = 1, 2$ or 3

Isothermal Euler equations + monofluid model

Isobaric pressure law:

$$\begin{aligned} p &= p_{\text{vapor}} = p_{\text{water}} \\ &= p_0 + c^2(\rho - (\alpha(\rho, \varphi)\rho_A + (1 - \alpha(\rho, \varphi))\rho_W)) \end{aligned} \quad (4)$$

where

$$\alpha(\rho, \varphi) = \frac{\theta - 1 + \sqrt{(\theta - 1)^2 + 4\theta\varphi}}{2\theta}$$

and

$$\theta = \frac{\rho_W - \rho_A}{\rho}$$

Volume fraction of gas: $\alpha(\rho, \varphi)$

Reference pressure: $p_0 = 10^5 Pa$

Reference density for water: $\rho_W = 1000 kg/m^3$

Reference density for air: $\rho_A = 1 kg/m^3$

Sound of speed: $c = 1500 m/s$

Boundary conditions

- $r \in [0, AR_b]$, R_b =initial bubble radius



$$Y = (\rho, u, p, \varphi),$$
$$Y(r, 0) = \begin{cases} Y_L & \text{if } r < R_b, \\ Y_R & \text{if } r > R_b. \end{cases} \quad (5)$$

- $Y(0, t) = Y_L$ and $Y(AR_b, t) = Y_R$
- $d > 1 \Rightarrow$ left boundary useless (because the first edge surface is 0)

Initial data

Reference parameters:

- Sound speed: $c_0 = 1500 m/s$
- Pressure: $p_0 = 10^5 Pa$
- Density: $\rho_0 = 1000 kg/m^3$

Vapor: $\gamma_1 = 1.4$ ($\pi_1 = 0$)

Water: $\gamma_2 = 1.1$, $\pi_2 = \frac{\rho_0 c_0^2}{\gamma_2}$

Left Side

$$\varphi_L = 1$$

$$u_L = 0$$

$$p_L = p_0 \left(\frac{R_{eq}}{R_b} \right)^{3\gamma_1} = 75208.8 Pa$$

$$\rho_L = \rho_{eq} \left(\frac{p_L}{p_0} \right)^{1/\gamma_1} = 0.948 kg/m^3$$

Right Side

$$\varphi_R = 0$$

$$u_R = 0$$

$$p_R = p_0$$

$$\rho_R = \rho_0$$

Plan

- 1 Studied models
 - Euler System
 - Isothermal Euler equations
 - Boundary conditions and initial data
- 2 Numerical scheme
- 3 Numerical results

Numerical scheme

Eulerian approach: poor precision, smears the interface

Lagrangian approach: more precise at the interface,
constant states (u, p) when $d=1$,

if $R_b \rightarrow 0 \Rightarrow \Delta t \rightarrow 0$
(because of the CFL condition)

\Rightarrow Lagrangian approach only at the interface

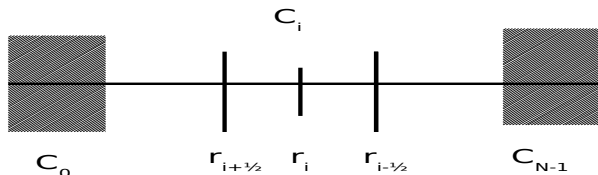
Numerical scheme

Conservative variables:

$$W = \begin{pmatrix} \rho\varphi \\ \rho(1-\varphi) \\ \rho u \\ \rho E \end{pmatrix} \quad \text{and} \quad W_i^n = W(r_i^n, t^n)$$

Left boundary motionless

Time step: $\tau^n = t^{n+1} - t^n$



Numerical scheme

$$r_{i+1/2}^{n+1} = r_{i+1/2}^n + \tau^n u_{i+1/2}^n. \quad (6)$$

$u_{i+1/2}^n = 0 \quad \longrightarrow \quad$ Eulerian scheme

$u_{i+1/2}^n = v_{i+1/2}^n \quad \longrightarrow \quad$ Lagrangian scheme

Numerical scheme

Integrating the balance laws on the space-time trapezoid

$$\{(r_{i-1/2}^n, t^n), (r_{i+1/2}^n, t^n), (r_{i+1/2}^{n+1}, t^{n+1}), (r_{i-1/2}^{n+1}, t^{n+1})\},$$

$$V_i^{n+1} W_i^{n+1} - V_i^n W_i^n + \tau^n \left(S_{i+1/2}^n F_{i+1/2}^n - S_{i-1/2}^n F_{i-1/2}^n \right) = \tau^n G_i^n. \quad (7)$$

Volume of ce cell i : $V_i^n = \int_{r_{i-1/2}^n}^{r_{i+1/2}^n} r^{d-1} dr$

Surface of edge : $i + 1/2$: $S_{i+1/2}^n = \left(r_{i+1/2}^n \right)^{d-1}$

Lagrangian flux : $F_{i+1/2}^n$

Dimensional source term : G_i^n

Numerical scheme

- 1 Solve de Riemann problem $R(W_L, W_R, x/t) = W(x, t)$

$$\begin{aligned} W_t + f(W)_x &= 0, \\ W(0, t) &= \begin{cases} W_L & \text{if } x < 0, \\ W_R & \text{if } x > 0. \end{cases} \end{aligned} \quad (8)$$

To take into account the edges velocity, we compute:

$$W_{i+1/2}^n = R(W_i^n, W_{i+1}^n, u_{i+1/2}^n)$$

2 Compute the Lagrangian flux

$$F_{i+1/2}^n = f(W_{i+1/2}^n) - v_{i+1/2}^n W_{i+1/2}^n$$

3 Compute the dimensional source term :

$$G_i^n \simeq \int_{r_{i-1/2}^n}^{r_{i+1/2}^n} (d-1)p(r)r^{d-2}dr$$

given by $G_i^n = p_i^n \left(S_{i+1/2}^n - S_{i-1/2}^n \right)$

Stability condition

$$h_i^n = \frac{V_i^n}{S_{i-1/2}^n + S_{i+1/2}^n}$$

The time step has to respect two conditions :

Non vanishing condition:

$$\tau^n \leq \max_i \left(\frac{h_i^n}{|v_{i\pm 1/2}^n|} \right) \quad (9)$$

Stability condition :

$$\tau^n \leq \max_i \left(\frac{h_i^n}{|u_{i\pm 1/2}^n - v_{i\pm 1/2}^n| + c_{i\pm 1/2}^n} \right) \quad (10)$$

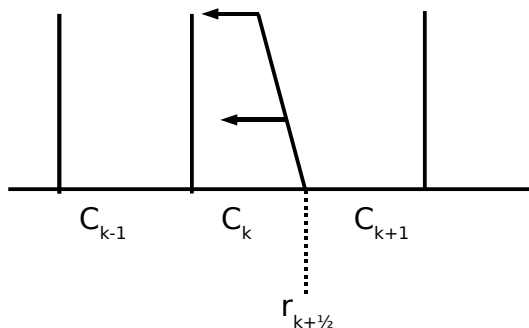
Cell averaging

Only the bubble interface move

$$\begin{aligned} \varphi_k^n \neq \varphi_{k+1}^n &\Rightarrow v_{k+1/2}^n = u_{k+1/2}^n, \\ v_{i+1/2}^n &= 0 \quad \text{if } i \neq k. \end{aligned} \tag{11}$$

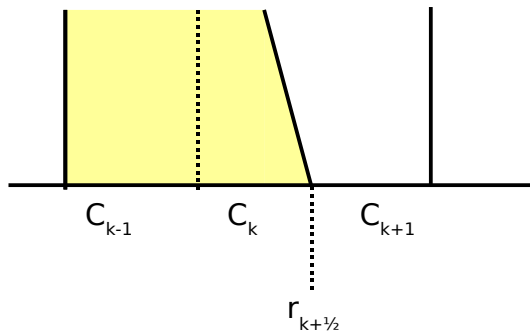
Cell averaging

The interface moves from right to left



Cell averaging

The interface moves from right to left



Cell averaging

- 1 Gather the two left cells
- 2 $\tau^n = \min(\tau^n, \Delta t')$ where

$$\Delta t' = \frac{r_{k+1/2}^n - r_{k-1/2}^n}{|v_{k+1/2}^n|}. \quad (12)$$

In the isothermal case

- 1 φ constant \Rightarrow Eulerian scheme
- 2 Godunov scheme is not employed (because the exact Riemann solver is too much CPU consuming) \Rightarrow Rusanov scheme

$$f_{i+\frac{1}{2}}^n = \frac{f_i^n + f_{i+1}^n}{2} - \frac{\sigma_{i+\frac{1}{2}}^n}{2} (w_{i+1}^n - w_i^n)$$

Plan

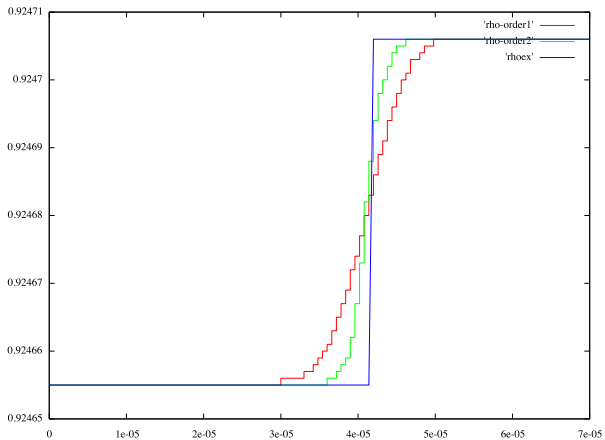
- 1 Studied models
 - Euler System
 - Isothermal Euler equations
 - Boundary conditions and initial data
- 2 Numerical scheme
- 3 Numerical results

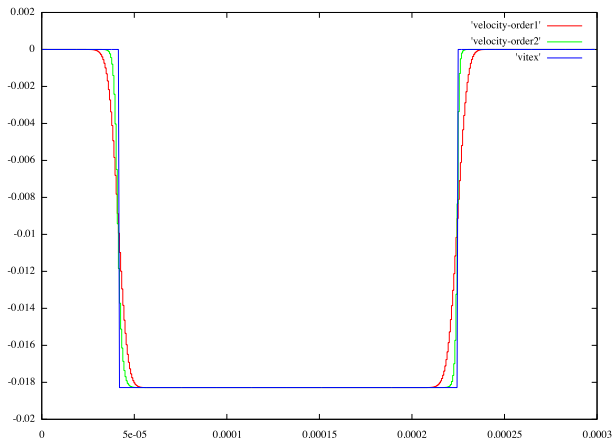
Validate ALE + averaging approach in 1D

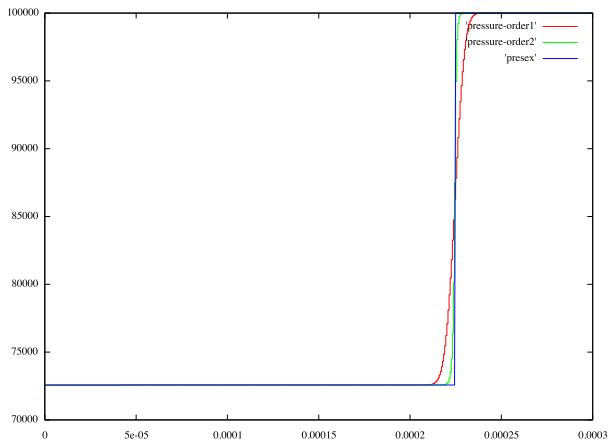
500 cells, CFL=0.8

Initial bubble radius $R_b = 0.7469 \times 10^{-4}$, $t_{final} = 1 \times 10^{-7} s$,
 $r \in [0, 4 * R_b]$

Comparison between: Exact Riemann Solver and Godunov with
ALE and stiffened gas law







Easy bubble test with $d=3$

$c_0 = 50m/s$: decrease the numerical viscosity of the scheme

$R_b = 0.7469 \times 10^{-4}$: oscillations of smaller amplitude

Left Side

$$\varphi_L = 1$$

$$u_L = 0$$

$$p_L = p_0 \left(\frac{R_{eq}}{R_b} \right)^{3\gamma_1}$$

$$\rho_L = \rho_{eq} \left(\frac{p_L}{p_0} \right)^{1/\gamma_1}$$

Right Side

$$\varphi_R = 0$$

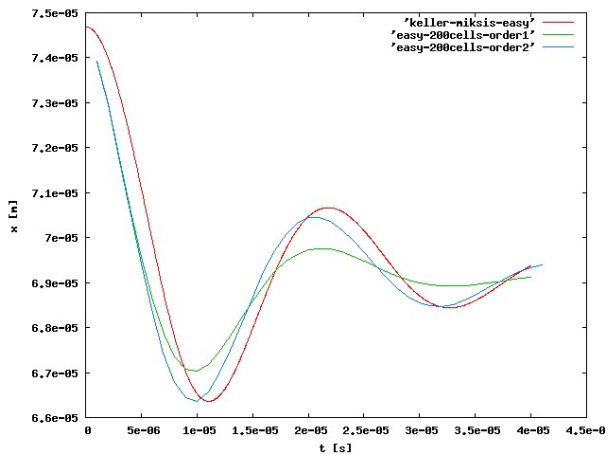
$$u_R = 0$$

$$p_R = P_0$$

$$\rho_R = \rho_0$$

$N=100$ cells, $CFL=0.8$, $t_{final} = 1 \times 10^{-7}s$.

Comparison between: Keller-Miksis model and Godunov with ALE and stiffened gas law



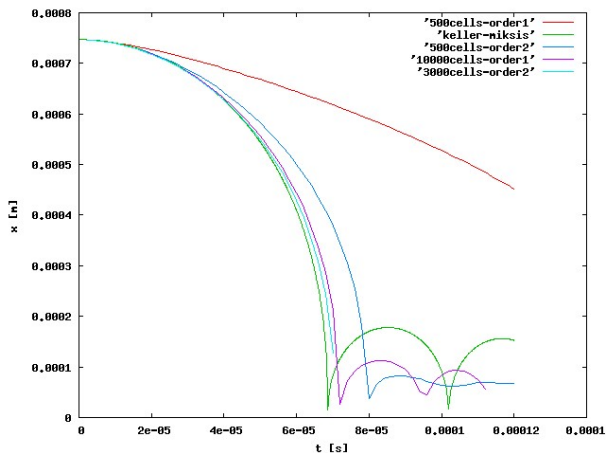
Hard bubble test with $d=3$

$c_0 = 1500m/s$: decrease the numerical viscosity of the scheme

$R_b = 0.7469 \times 10^{-3}$: oscillations of smaller amplitude

3000 and 10000 cells

Comparison between: Keller-Miksis model and Godunov with ALE
and stiffened gas law

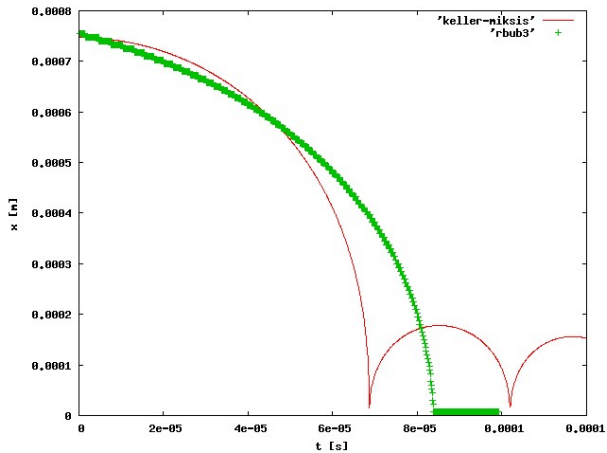


Isothermal case with $d=3$

2000 cells

$$\varphi = 10^{-3}$$

Comparison between: Keller-Miksis model and Rusanov with isobaric pressure law





Conclusion


Future works

- 1 Second order in time
- 2 Employ the VFRoe scheme
- 3 Isothermal two-fluid model with “linear” pressure law

References

 Saurel, Richard and Abgrall, Rémi
A simple method for compressible multifluid flows.
SIAM Journal on Scientific Computing, 1999.

 Hirt, C. W. and Amsden, A. A. and Cook, J. L.
An arbitrary Lagrangian-Eulerian computing method for all
flow speeds
Journal of Computational Physics, 1997.

 Barberon, T. and Helluy, P. and Rouy, S.
Practical computation of axisymmetrical multifluid flows
International Journal of Finite Volumes, 2003.

References



Godlewski, Edwige and Raviart, Pierre-Arnaud
Numerical approximation of hyperbolic systems of conservation laws

Applied Mathematical Sciences, Springer-Verlag, 1996.



Helluy, P. and Seguin, N.
Relaxation models of phase transition flows

Mathematical Modelling and Numerical Analysis, 2005.



Keller, J. B. and Miksis, M.
Bubble oscillations of large amplitude

Acoustical Society of America Journal, 1980.