TU München Zentrum Mathematik Lehrstuhl M16 E. Sonnendrücker E. Franck

Numerical methods for hyperbolic systems

Exercise sheet 3: Linear hyperbolic systems

Exercice 1

We consider the wave equation

$$\begin{cases} \partial_t p + \partial_x u = 0, \\ \partial_t u + \partial_x p = 0, \end{cases}$$
(1)

1. Diagonalize the system and show that the upwind scheme is given by

$$\begin{cases} \frac{v_j^{n+1} - v_j^n}{\Delta t} + \frac{v_j^n - v_{j-1}^n}{\Delta x} = 0, \\ \frac{w_j^{n+1} - w_j^n}{\Delta t} - \frac{w_{j+1}^n - w_j^n}{\Delta x} = 0, \end{cases}$$
(2)

with v = p + u and w = p - u.

2. Prove that the upwind scheme (2) for the initial system (1) can be write on the following form

$$\begin{cases} \frac{p_j^{n+1} - p_j^n}{\Delta t} + \frac{u_{j+1}^n - u_{j-1}^n}{2\Delta x} - \frac{p_{j+1}^n - 2p_j^n + p_{j-1}^n}{2\Delta x} = 0, \\ \frac{u_j^{n+1} - u_j^n}{\Delta t} + \frac{p_{j+1}^n - p_{j-1}^n}{2\Delta x} - \frac{u_{j+1}^n - 2u_j^n + u_{j-1}^n}{2\Delta x} = 0, \end{cases}$$
(3)

and used this form to compute the consistency error.

3. Prove that the scheme (2) satisfy the maximum principle under a CFL condition.

4. Prove that the scheme (2) is stable for all l^q norms $(1 \le q \le \infty)$ using the previous result and convex functions. The l^q norm is defined by

$$||(v,w)||_{l^q} = \left(\Delta x \sum_j ||(v_j,w_j)||_q^q\right)^{\frac{1}{q}}$$

with $||(v_j, w_j)||_q^q = |v_j|^q + |w_j|^q$.

Additional question

We introduce the damped wave equation

$$\begin{cases} \partial_t p + \partial_x u = 0, \\ \partial_t u + \partial_x p = -\sigma u, \end{cases}$$
(4)

and the upwind scheme associated

$$\begin{cases} \frac{p_j^{n+1} - p_j^n}{\Delta t} + \frac{u_{j+1}^n - u_{j-1}^n}{2\Delta x} - \frac{p_{j+1}^n - 2p_j^n + p_{j-1}^n}{2\Delta x} = 0, \\ \frac{u_j^{n+1} - u_j^n}{\Delta t} + \frac{p_{j+1}^n - p_{j-1}^n}{2\Delta x} - \frac{u_{j+1}^n - 2u_j^n + u_{j-1}^n}{2\Delta x} = -\sigma u_j^n, \end{cases}$$
(5)

5. We call "steady states" the solutions of the systems defined by $\partial_x u = 0$ and $\partial_x p = -\sigma u$. Prove that (5) preserve exactly the steady states.

Exercice 2

We consider the Maxwell equation

$$\begin{cases} \mu \partial_t B + c \partial_x E = -c \sigma^* B, \\ \varepsilon \partial_t E + c \partial_x B = -c \sigma E, \end{cases}$$
(6)

with periodic boundary condition on $\Omega = [0, L]$ and $\mu > 0$, $\varepsilon > 0$, $\sigma > 0$, $\sigma^* > 0$, c > 0.

1. Prove the following energy equality and the uniqueness of the solutions.

$$\frac{d}{dt}\left(\int_{\Omega}\varepsilon|E(t,x)|^2+\mu|B(t,x)|^2dx\right) = -\int_{\Omega}\sigma|E(t,x)|^2+\sigma^*\mu|B(t,x)|^2dx\tag{7}$$

2. We introduce the plane waves (which are a good approximations of physical waves) defined by $E(t,x) = E_0 e^{j(wt-kx)}$ and $B(t,x) = B_0 e^{j(wt-kx)}$ (j complex number) with $E_0 \in \mathbb{R}$, $B_0 \in \mathbb{R}$, k the wave vector and w the frequency. Give the conditions (called dispersion relation) on w and k such as the plane waves are solutions of (6) for $\sigma = 0$ and $\sigma^* = 0$.

In this part we assume that $\sigma = 0$ and $\sigma^* = 0$. Now we introduce the DG centered scheme for (6). The mesh Ω_h is defined by n + 1 points x_i and n cells $K_i = \left[x_{i-\frac{1}{2}}, x_{i+\frac{1}{2}}\right]$. The volume of the cell K_i is $\Delta x_i = |x_{i+\frac{1}{2}} - x_{i-\frac{1}{2}}|$. We call a generic cell K. To finish the test function are defined by $v \in V_h = \{v/v|_K \in \mathbb{P}^p(K)\}$. The scheme is given by

$$\begin{cases} \varepsilon \sum_{l=0}^{k} \int_{K_{i}} \phi_{l}^{i} \phi_{m}^{i} \left(\frac{E_{l,i}^{n+1} - E_{l,i}^{n}}{\Delta t} \right) - c \sum_{l=0}^{k} B_{l,i}^{n} \int_{K_{i}} \phi_{l}^{i} \partial_{x} \phi_{m}^{i} + c \sum_{l=0}^{k} \left[B \phi_{m}^{i} \right]_{x_{i-\frac{1}{2}}}^{x_{i+\frac{1}{2}}} = 0, \quad \forall 0 \le m \le k, \\ \mu \sum_{l=0}^{k} \int_{K_{i}} \phi_{l}^{i} \phi_{m}^{i} \left(\frac{B_{l,i}^{n+1} - B_{l,i}^{n}}{\Delta t} \right) - c \sum_{l=0}^{k} E_{l,i}^{n} \int_{K_{i}} \phi_{l}^{i} \partial_{x} \phi_{m}^{i} + c \sum_{l=0}^{k} \left[E \phi_{m}^{i} \right]_{x_{i-\frac{1}{2}}}^{x_{i+\frac{1}{2}}} = 0, \quad \forall 0 \le m \le k, \end{cases}$$

$$\tag{8}$$

with $\left[B\phi_m^i\right]_{x_{i-\frac{1}{2}}}^{x_{i+\frac{1}{2}}} = \frac{1}{2}\left(B_{l,i+1}^n\phi_l^{i+1}(x_{i+\frac{1}{2}})\phi_m^i(x_{i+\frac{1}{2}}) + B_{l,i}^n\phi_l^i(x_{i+\frac{1}{2}})\phi_m^i(x_{i+\frac{1}{2}})\right)$

$$-\frac{1}{2}\left(B_{l,i-1}^{n}\phi_{l}^{i-1}(x_{i-\frac{1}{2}})\phi_{m}^{i}(x_{i-\frac{1}{2}})+B_{l,i}^{n}\phi_{l}^{i}(x_{i-\frac{1}{2}})\phi_{m}^{i}(x_{i-\frac{1}{2}})\right).$$

3. We consider $V_h = P^1(K)$. We propose to use the Lagrange polynomial associated with the point $x_{i-\frac{1}{2}}$ and $x_{i+\frac{1}{2}}$. Prove that the family is a basis of V_h . Write the scheme in a cell K_i .

4. In this exercise we propose to study the numerical dispersion relation which define the numerical wave vector \tilde{k} for $V_h = P^0(K) = \text{Span}(1)$. Write the scheme for this basis. Now we define $B_i^n = B_0 e^{j(wn\Delta t - \tilde{k}i\Delta x)}$ and $E_i^n = E_0 e^{j(wn\Delta t - \tilde{k}i\Delta x)}$ with j the complex number and i the index of the cell. Gives the relation between w and \tilde{k} such as the discrete plane waves are solutions of (8). Show that the numerical dispersive relation is $\tilde{k}^2 = \frac{w^2}{c^2} + O(\Delta x^p + \Delta t^q)$ with p > 1 and q > 1.