

# A fully well-balanced hydrodynamic reconstruction

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## Problem statement

The hydrodynamic reconstruction

Suitable expression of  $\mathcal{H}$

Linear high-order extension

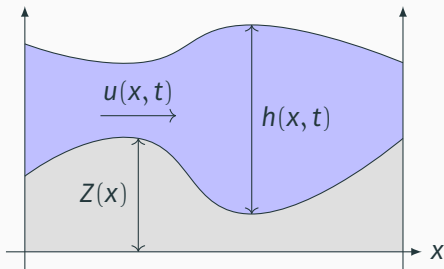
Numerical experiments

Conclusion

# The shallow water equations with topography

$$\begin{cases} \partial_t h + \partial_x q = 0, \\ \partial_t q + \partial_x \left( \frac{q^2}{h} + \frac{1}{2} g h^2 \right) = -g h \partial_x Z(x) \end{cases}$$

The equations are written under the form  $\partial_t W + \partial_x F(W) = S(W)$ .



- $h(x, t)$ : water height
- $u(x, t)$ : water velocity
- $q = hu$ : water discharge
- $Z(x)$ : known topography
- $g$ : gravity constant

We pay particular attention to solutions of prime importance:  
the **steady solutions**.

## Shallow water with topography: steady solutions

Taking  $\partial_t W = 0$  in the shallow water system yields

$$\left\{ \begin{array}{l} \partial_x q = 0, \\ \partial_x \left( \frac{q^2}{h} + \frac{1}{2} g h^2 \right) = -g h \partial_x Z, \end{array} \right. \xrightarrow[\text{solution}]{\text{smooth}} \left\{ \begin{array}{l} q = \text{cst} = q_0, \\ \partial_x \left( \frac{q_0^2}{2h^2} + g(h + Z) \right) = 0. \end{array} \right.$$

We summarize the second relation by introducing a function  $B$  such that, for a steady solution,  $B(h, q_0, Z) = B_0$ .

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**Two cases** are distinguished:

- $q_0 = 0 \rightsquigarrow$  lake at rest

we get  $B(h, q_0, Z) = h + Z = \text{cst}$ : linear equation in  $h$

- $q_0 \neq 0 \rightsquigarrow$  moving steady solution

we get  $B(h, q_0, Z) = \frac{q_0^2}{2h^2} + g(h + Z) = B_0$ : **nonlinear** equation in  $h$ !

# Finite volume scheme

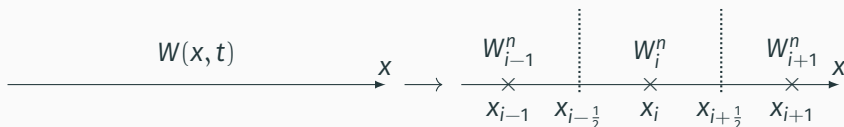
Recall the compact form of the shallow water equations:

$$\partial_t W + \partial_x F(W) = S(W).$$

We take a **generic finite volume numerical scheme** approximating the shallow water equations:

$$\frac{W_i^{n+1} - W_i^n}{\Delta t} + \frac{1}{\Delta x} \left[ \mathcal{F}(W_i^n, W_{i+1}^n) - \mathcal{F}(W_{i-1}^n, W_i^n) \right] = \mathcal{S}(W_{i-1}^n, W_i^n, W_{i+1}^n),$$

with  $\mathcal{F}$  a **consistent numerical flux**, i.e.  $\mathcal{F}(W, W) = F(W)$ , and  $\mathcal{S}$  a consistent numerical source term.



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with  $\mathcal{F}$  a **consistent numerical flux**, i.e.  $\mathcal{F}(W, W) = F(W)$ , and  $\mathcal{S}$  a consistent numerical source term.

## Definition: well-balanced scheme

A numerical method approximating the solution of a balance law is called **well-balanced** if it exactly preserves the steady solutions.

**Question:** can we make this **generic finite volume scheme** **well-balanced** without changing the numerical flux?

# An answer for the lake at rest: the hydrostatic reconstruction

The **hydrostatic reconstruction** was introduced in E. Audusse et al., *SIAM J. Sci. Comput.* (2004), as a way to make it possible for any finite volume scheme to capture the **lake at rest** steady solution.

It relies on:

1. providing a relevant expression for  $\mathcal{S}$ ,
2. evaluating the numerical flux at a specific reconstruction of  $W$ .

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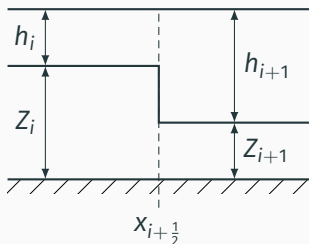
$$\frac{W_i^{n+1} - W_i^n}{\Delta t} + \frac{1}{\Delta x} \left[ \mathcal{F}(W_{i+\frac{1}{2},-}^n, W_{i+\frac{1}{2},+}^n) - \mathcal{F}(W_{i-\frac{1}{2},-}^n, W_{i-\frac{1}{2},+}^n) \right] = \mathcal{S}_i^n$$

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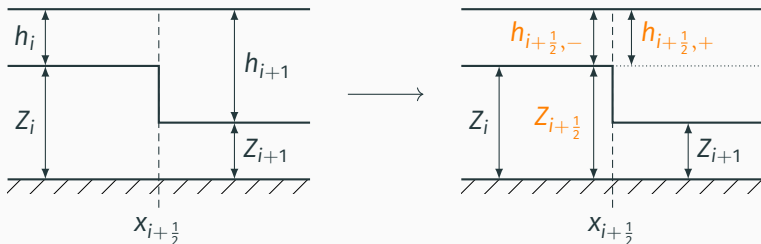


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It relies on:

1. providing a relevant expression for  $\mathcal{S}$ ,
2. evaluating the **numerical flux** at a **specific reconstruction** of  $W$ .



**Main goal of this work:** Provide a **linear** reconstruction able to capture the steady solutions with  $q_0 = 0$  or  $q_0 \neq 0$ .

The objectives of our **hydrodynamic reconstruction** include:

- making sure that the resulting scheme is **consistent**,
- ensuring the **capture of steady solutions** with  $q_0 = 0$  or  $q_0 \neq 0$ ,
- handling **dry areas** and transitions between wet and dry areas (not presented in this talk),
- a **linear and well-balanced high-order extension**.

Problem statement

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Suitable expression of  $\mathcal{H}$

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# Expression of the hydrodynamic reconstruction

Away from dry areas, the **hydrostatic reconstruction** reads:

$$h_{i+\frac{1}{2},-}^n = h_i^n + \left( z_i - z_{i+\frac{1}{2}} \right),$$

$$h_{i+\frac{1}{2},+}^n = h_{i+1}^n + \left( z_{i+1} - z_{i+\frac{1}{2}} \right).$$

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with  $\mathcal{H}$  a function of  $h_L$ ,  $h_R$ ,  $\bar{q}$  and  $\Delta Z := Z_R - Z_L$ , and with

$$\text{Fr}^2(h_L, h_R, \bar{q}) = \frac{\bar{q}^2(h_L + h_R)}{2gh_L^2h_R^2}.$$

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**The hydrodynamic reconstruction relies on deriving a suitable function  $\mathcal{H}$ .**



# Characterization of interface steady relations

$$h_{i+\frac{1}{2},-}^n = h_i^n + \left( Z_i - Z_{i+\frac{1}{2}} \right) + 2\text{Fr}^2 \left( h_i^n, h_{i+\frac{1}{2}}^n, q_i^n \right) \mathcal{H} \left( h_i^n, h_{i+\frac{1}{2}}^n, q_i^n, Z_{i+\frac{1}{2}} - Z_i \right)$$

Define the interface state by

$$(h_{i+\frac{1}{2}}^n, Z_{i+\frac{1}{2}}) = \begin{cases} (h_i^n, Z_i) & \text{if } Z_i > Z_{i+1}, \\ (h_{i+1}^n, Z_{i+1}) & \text{otherwise.} \end{cases}$$

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When the solution is steady, setting  $\bar{q} = q_i = q_{i+1}$ , we get:

$$B(h_i, \bar{q}, Z_i) = B(h_{i+1}, \bar{q}, Z_{i+1}).$$

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## Well-balancing requirement on $\mathcal{H}$

Some algebraic manipulations allow us to write

$$\frac{\bar{q}^2}{2h_i^2} + g(h_i + Z_i) = \frac{\bar{q}^2}{2h_{i+\frac{1}{2}}^2} + g(h_{i+\frac{1}{2}} + Z_{i+\frac{1}{2}})$$

$$\Longleftrightarrow$$

$$Z_{i+\frac{1}{2}} - Z_i = -\left(h_{i+\frac{1}{2}} - h_i\right) \left(1 - \text{Fr}^2\left(h_i, h_{i+\frac{1}{2}}, \bar{q}\right)\right),$$

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which is nothing but the usual discrete characterization of smooth steady solutions.

We claim that imposing the following property on  $\mathcal{H}$  will be enough to preserve steady solutions:

$$\Delta Z = -(h_R - h_L)(1 - \text{Fr}^2(h_L, h_R, \bar{q})) \implies \mathcal{H}(h_L, h_R, \bar{q}, \Delta Z) = \frac{h_R - h_L}{2}.$$

## Well-balancing requirement on $\mathcal{H}$

Indeed, assuming that the solution is steady, we obtain the following sequence of equalities:

$$\begin{aligned} h_{i+\frac{1}{2},-}^n &= h_i^n + \left( Z_i - Z_{i+\frac{1}{2}} \right) \\ &\quad + 2\text{Fr}^2 \left( h_i^n, h_{i+\frac{1}{2}}^n, q_i^n \right) \mathcal{H} \left( h_i^n, h_{i+\frac{1}{2}}^n, q_i^n, Z_{i+\frac{1}{2}} - Z_i \right), \end{aligned}$$

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$$h_{i+\frac{1}{2},-}^n = h_{i+\frac{1}{2}}^n,$$

which proves that the scheme is well-balanced.

## Summary and source term discretization

To summarize, for the reconstruction to be **consistent** and **well-balanced**, we require the **following two properties** on the bounded function  $\mathcal{H}$ :

1.  $\mathcal{H}(h_L, h_R, \bar{q}, \Delta Z) \underset{\Delta Z \rightarrow 0}{=} \mathcal{O}(\Delta Z),$
2.  $\Delta Z = -(h_R - h_L)(1 - \text{Fr}^2(h_L, h_R, \bar{q})) \implies \mathcal{H}(h_L, h_R, \bar{q}, \Delta Z) = \frac{h_R - h_L}{2}.$

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In addition, the whole scheme will also be consistent and well-balanced if the following **numerical source term** is used:

$$\Delta x(\mathcal{S}_q)_i^n = -g \frac{2h_{i-\frac{1}{2},+}^n h_{i+\frac{1}{2},-}^n}{h_{i-\frac{1}{2},+}^n + h_{i+\frac{1}{2},-}^n} (Z_{i+\frac{1}{2}} - Z_{i-\frac{1}{2}}) + \frac{4g}{h_{i-\frac{1}{2},+}^n + h_{i+\frac{1}{2},-}^n} \mathcal{H}(h_{i-\frac{1}{2},+}^n, h_{i+\frac{1}{2},-}^n, q_i, Z_{i+\frac{1}{2}} - Z_{i-\frac{1}{2}})^3.$$

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The proof results from algebraic manipulations (not detailed here).

**Next step: obtain a suitable expression of  $\mathcal{H}$ .**

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**Suitable expression of  $\mathcal{H}$**

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# Satisfying the well-balanced property

Recall that we need

$$\mathcal{H}(h_L, h_R, \bar{q}, \Delta Z) = \frac{h_R - h_L}{2}$$

when  $\mathcal{H}$  is applied to a discrete steady solution.

To obtain an expression of  $\mathcal{H}$  satisfying this property, we need to understand how  $(h_R - h_L)/2$  behaves for discrete steady solutions.

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To obtain an expression of  $\mathcal{H}$  satisfying this property, we need to **understand how  $(h_R - h_L)/2$  behaves for discrete steady solutions.**

We now seek a relation to characterize **the jump of  $h$  at the interface**, i.e. an expression of  $(h_R - h_L)/2$  for steady solutions.

We **assume that the solution is steady**, and introduce notation

$$\bar{h} := \frac{h_L + h_R}{2} \quad \text{and} \quad [h] := \frac{h_R - h_L}{2},$$

so that  $h_L$  and  $h_R$  satisfy

$$h_L = \bar{h} - [h] \quad \text{and} \quad h_R = \bar{h} + [h].$$

The goal is now to rewrite the steady relation in terms of  $\bar{h}$  and  $[h]$ .



# A local relation to characterize steady solutions

Recall that the **steady solutions** are governed by

$$B(h, q_0, Z) = \frac{q_0^2}{2h^2} + g(h + Z) = B_0.$$

That is to say, at the interface between states  $W_L$  and  $W_R$ , the solution is locally steady if  $q_L = q_R = \bar{q}$  and

$$B(h_L, \bar{q}, Z_L) = B(h_R, \bar{q}, Z_R) \iff \frac{\bar{q}^2}{2h_L^2} + g(h_L + Z_L) = \frac{\bar{q}^2}{2h_R^2} + g(h_R + Z_R).$$

We set out to rewrite **the above relation** using  $\bar{h}$  and  $[h]$  instead of  $h_L$  and  $h_R$ .

## A nonlinear relation for the interface jump

$$\frac{\bar{q}^2}{2h_L^2} + g(h_L + Z_L) = \frac{\bar{q}^2}{2h_R^2} + g(h_R + Z_R)$$

$$\Longleftrightarrow$$

$$\frac{\bar{q}^2}{2(\bar{h} - [h])^2} + g(\bar{h} - [h] + Z_L) = \frac{\bar{q}^2}{2(\bar{h} + [h])^2} + g(\bar{h} + [h] + Z_R)$$

$$\Longleftrightarrow$$

$$\dots$$

$$\Longleftrightarrow$$

$$2[h] \left( g(\bar{h}^2 - [h]^2)^2 - \bar{q}^2 \bar{h} \right) = -g(Z_R - Z_L) (\bar{h}^2 - [h]^2)^2.$$

## “Quadratized” relation

$$2\mathcal{H}\left(g(\bar{h}^2 - \mathcal{H}^2)^2 - \bar{q}^2\bar{h}\right) = -g\Delta Z(\bar{h}^2 - \mathcal{H}^2)^2 \quad (*)$$

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Equation (\*) is nonlinear, and using it would incur considerable computational cost. To avoid this issue, we proceed with a **linearization**-like simplification. First, for  $\mathcal{H} \neq \bar{h}$ , we get

$$(*) \iff 2\mathcal{H}\left(1 - \frac{\bar{q}^2\bar{h}}{g(\bar{h}^2 - \mathcal{H}^2)^2}\right) = -\Delta Z.$$

We then choose a **“quadratization”** of this expression around  $\mathcal{H} = [h]$ :

$$2\mathcal{H}\left(1 - \frac{\bar{q}^2(h_L + h_R)}{2g(\bar{h}^2 - [h]^2)} + \alpha([h] - \mathcal{H})\right) = -\Delta Z.$$

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In practice, after some testing, we choose

$$a = \text{sgn}(\Delta Z) \sqrt{\frac{|\Delta Z|}{2|[h]^3}}.$$

## Final expression of $\mathcal{H}$

We are left with  $\mathcal{H}$  satisfying a quadratic relation.

Solving this quadratic equation for  $\mathcal{H}$  leads to

$$\mathcal{H} = \frac{1}{2} \left( E - \operatorname{sgn}(1 - \operatorname{Fr}^2) \operatorname{sgn}(\Delta Z) \sqrt{E^2 + \sqrt{\frac{1}{2} |\Delta Z|} |[h]^3} \right),$$

with  $E = [h] + \frac{1 - \operatorname{Fr}^2}{2} \operatorname{sgn}(\Delta Z) \sqrt{\frac{|[h]^3}{2 |\Delta Z|}}.$

We show that, if  $\Delta Z$  and  $1 - \operatorname{Fr}^2$  do not simultaneously vanish:

1. this expression of  $\mathcal{H}$  is **well-balanced**;
2. this expression of  $\mathcal{H}$  is **consistent**, despite the divisions by  $\Delta Z$ .

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**Next step:** provide a well-balanced and linear high-order extension.



Problem statement

The hydrodynamic reconstruction

Suitable expression of  $\mathcal{H}$

**Linear high-order extension**

Numerical experiments

Conclusion

# Linear high-order extension

We follow the general strategy from [C. Berthon, S. Bulteau, F. Foucher, M. M'Baye and V. M.-D., *SIAM SISC*, 2022].

At each interface, we introduce a **convex combination of parameter**  $\theta_{i+\frac{1}{2}}$  between the **high-order reconstruction**  $W_{i+\frac{1}{2}}^{\text{HO}}$  and the **hydrodynamic reconstruction**  $W_{i+\frac{1}{2}}^{\text{HDR}}$ :

$$W_{i+\frac{1}{2},\pm} = \theta_{i+\frac{1}{2}} W_{i+\frac{1}{2},\pm}^{\text{HO}} + (1 - \theta_{i+\frac{1}{2}}) W_{i+\frac{1}{2},\pm}^{\text{HDR}}.$$

The coefficient  $\theta_{i+\frac{1}{2}}$  is based on the error to the steady solution, and

- if  $\theta_{i+\frac{1}{2}} = 0$ , the solution is steady, the scheme is well-balanced;
- if  $\theta_{i+\frac{1}{2}} = 1$ , the solution is unsteady, the scheme is high-order accurate.

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We provide several numerical tests with a finite volume scheme using the HLL flux:

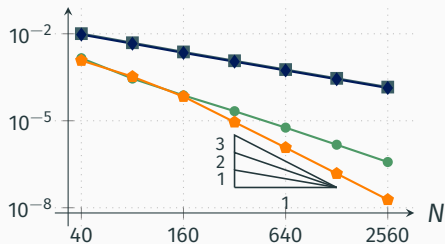
- an order of convergence test,
- three tests of the well-balanced property,
- a dam-break on a dry slope.

These tests are performed with the **h**ydrostatic **r**econstruction (HSR) and the **h**ydro**d**ynamic **r**econstruction (HDR).

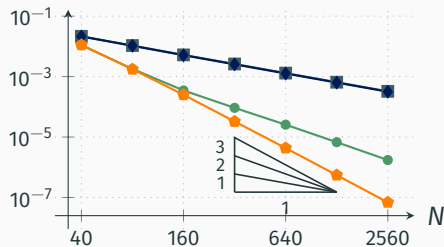
The schemes of order  $\delta$  are denoted by  $\text{HSR}\delta$  and  $\text{HDR}\delta$ .

# Order of convergence

$L^2$  error on  $h$

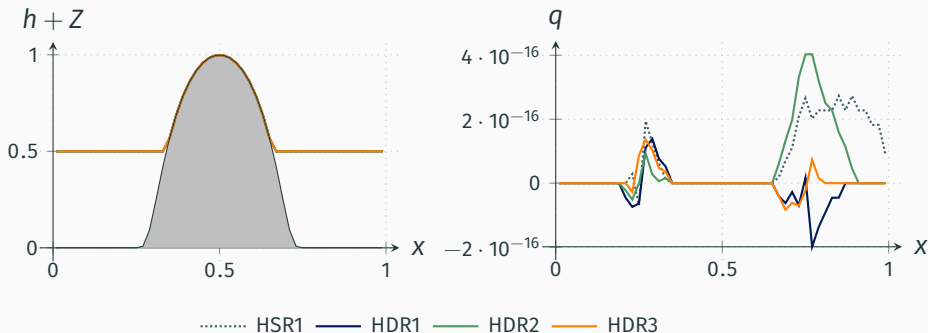


$L^2$  error on  $q$



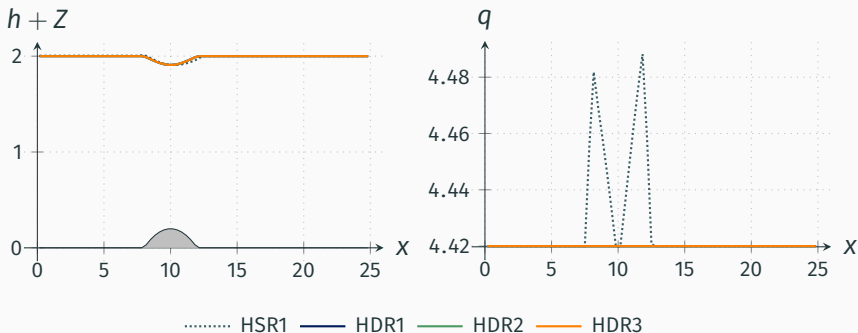
—■— HSR1    —◆— HDR1    —●— HDR2    —◆— HDR3

# Emerged lake at rest (50 cells)



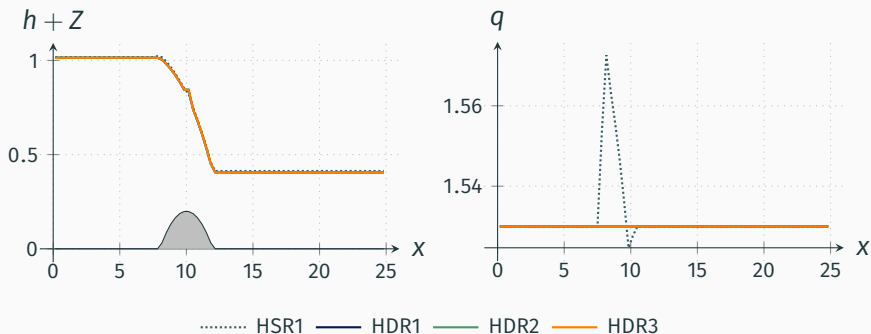
	HSR, $\mathbb{P}_0$	HDR, $\mathbb{P}_0$	HDR, $\mathbb{P}_1$	HDR, $\mathbb{P}_2$
$L^2$ error on $h$	$1.85 \cdot 10^{-17}$	$2.75 \cdot 10^{-17}$	$3.07 \cdot 10^{-17}$	$1.32 \cdot 10^{-17}$
$L^2$ error on $q$	$1.24 \cdot 10^{-16}$	$5.17 \cdot 10^{-17}$	$1.24 \cdot 10^{-16}$	$3.59 \cdot 10^{-17}$

# Subcritical flow (75 cells)



	HSR, $\mathbb{P}_0$	HDR, $\mathbb{P}_0$	HDR, $\mathbb{P}_1$	HDR, $\mathbb{P}_2$
$L^2$ error on $q$	$7.73 \cdot 10^{-2}$	$1.06 \cdot 10^{-14}$	$1.31 \cdot 10^{-14}$	$1.30 \cdot 10^{-14}$
$L^2$ error on $B$	$1.79 \cdot 10^{-1}$	$2.73 \cdot 10^{-14}$	$3.61 \cdot 10^{-14}$	$2.68 \cdot 10^{-14}$

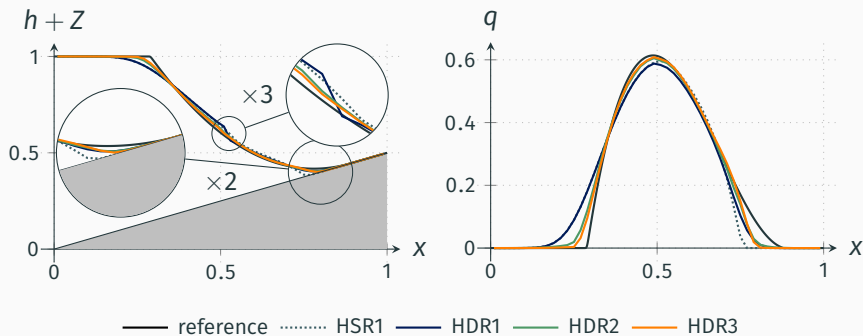
# Transcritical flow (75 cells)



	HSR, $\mathbb{P}_0$	HDR, $\mathbb{P}_0$	HDR, $\mathbb{P}_1$	HDR, $\mathbb{P}_2$
$L^2$ error on $q$	$3.74 \cdot 10^{-2}$	$4.73 \cdot 10^{-14}$	$5.15 \cdot 10^{-14}$	$5.21 \cdot 10^{-14}$
$L^2$ error on $B$	$1.45 \cdot 10^{-1}$	$4.50 \cdot 10^{-14}$	$5.12 \cdot 10^{-14}$	$5.92 \cdot 10^{-14}$



# Dam-break on a dry slope (50 cells)



Problem statement

The hydrodynamic reconstruction

Suitable expression of  $\mathcal{H}$

Linear high-order extension

Numerical experiments

**Conclusion**

We have developed a linear reconstruction that allows any finite volume scheme to be fully well-balanced for the shallow water system.

This reconstruction has the following properties:

- it leads to a **consistent** scheme,
- the resulting scheme is **well-balanced**,
- it is able to **handle wet/dry transitions**,
- it can be extended to **high-order accuracy** with no **nonlinear solver**.

This work led to the following preprint:

**C. Berthon, V. M.-D., *under review*, 2023.**

# Thank you for your attention!

Finite Volumes for Complex Applications 10 (**FVCA10**)

October 30, 2023 – November 03, 2023 in Strasbourg, France



## A nonlinear relation for the interface jump: properties

$$2\mathcal{H}\left(g(\bar{h}^2 - \mathcal{H}^2)^2 - \bar{q}^2\bar{h}\right) = -g\Delta Z(\bar{h}^2 - \mathcal{H}^2)^2 \quad (*)$$

Can  $\mathcal{H}$ , implicitly given by the above expression, satisfy the required consistency and well-balanced properties?

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Can  $\mathcal{H}$ , implicitly given by the above expression, satisfy the required **consistency** and **well-balanced** properties?

1. For the **consistency**, we need  $\mathcal{H}(h_L, h_R, \bar{q}, \Delta Z) \underset{\Delta Z \rightarrow 0}{=} \mathcal{O}(\Delta Z)$ : at least one solution to (\*) satisfies this property.

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2. For the **well-balanced property**, we need

$$\Delta Z = -(h_R - h_L)(1 - \text{Fr}^2(h_L, h_R, \bar{q})) \implies \mathcal{H}(h_L, h_R, \bar{q}, \Delta Z) = \frac{h_R - h_L}{2}.$$

This property holds since (\*) has been derived so that  $2\mathcal{H} = h_R - h_L$  is a solution as soon as the flow is steady.

## Well-balanced property

To show the **well-balanced** property, we take  $\Delta Z = -(1 - Fr^2)\Delta h$ , to get

$$E = \Delta h + \frac{1 - Fr^2}{4} \operatorname{sgn}(-(1 - Fr^2)\Delta h) \sqrt{\frac{|\Delta h|^3}{|1 - Fr^2||\Delta h|}} = \Delta h \left( 1 - \frac{1}{4} \sqrt{|1 - Fr^2|} \right),$$

$$E^2 + \sqrt{|\Delta Z||\Delta h|^3} = (\Delta h)^2 \left( 1 + \frac{1}{4} \sqrt{|1 - Fr^2|} \right)^2.$$



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$$E^2 + \sqrt{|\Delta Z||\Delta h|^3} = (\Delta h)^2 \left(1 + \frac{1}{4} \sqrt{|1 - Fr^2|}\right)^2.$$

Plugging this in  $\mathcal{H}$ , we obtain

$$\begin{aligned} \mathcal{H} &= \frac{1}{4} \left( \Delta h \left(1 - \frac{1}{4} \sqrt{|1 - Fr^2|}\right) + \operatorname{sgn}(\Delta h) \sqrt{(\Delta h)^2 \left(1 + \frac{1}{4} \sqrt{|1 - Fr^2|}\right)^2} \right) \\ &= \frac{\Delta h}{4} \left( 1 - \frac{1}{4} \sqrt{|1 - Fr^2|} + 1 + \frac{1}{4} \sqrt{|1 - Fr^2|} \right) = \frac{\Delta h}{2}, \end{aligned}$$

which proves the well-balanced property.

# High-order scheme

A high-order (non-well-balanced) finite volume scheme reads:

$$W_i^{n+1} = W_i^n - \frac{\Delta t}{\Delta x} \left( \mathcal{F}(\widehat{W}_{i,+}^n, \widehat{W}_{i+1,-}^n) - \mathcal{F}(\widehat{W}_{i-1,+}^n, \widehat{W}_{i,-}^n) \right) + \Delta t \widehat{S}_i^n.$$

In each cell, we reconstruct a polynomial of degree  $d$ , under the form

$$\widehat{W}_i^n(x) = W_i^n + \sum_{\alpha=1}^d R_i^\alpha (x - x_i)^\alpha,$$

where the coefficients  $R_i^\alpha$  depend on the neighboring cells.

The evaluations at the interfaces  $x_{i \pm \frac{1}{2}}$  are then given by:

$$\widehat{W}_{i,-}^n = W_i^n + \sum_{\alpha=1}^d R_i^\alpha \left( -\frac{\Delta x}{2} \right)^\alpha \quad \text{and} \quad \widehat{W}_{i,+}^n = W_i^n + \sum_{\alpha=1}^d R_i^\alpha \left( \frac{\Delta x}{2} \right)^\alpha,$$

and the high-order source term is the following approximation:

$$\widehat{S}_i^n = \frac{1}{\Delta x} \int_{x_{i-\frac{1}{2}}}^{x_{i+\frac{1}{2}}} S(\widehat{W}_i^n(x)) dx + \mathcal{O}(\Delta x^{d+1}).$$

# Linear well-balanced correction of the high-order scheme

We introduce a **convex combination with parameter  $\theta_{i\pm\frac{1}{2}}$**  to provide a **well-balanced** correction to the **high-order scheme**, such that:

- if  $\theta_{i\pm\frac{1}{2}} = 0$ , the scheme is well-balanced;
- if  $\theta_{i\pm\frac{1}{2}} = 1$ , the scheme is high-order accurate.

The new evaluations at the interfaces  $x_{i\pm\frac{1}{2}}$  are given by:

$$\tilde{W}_{i,-}^n = W_i^n + \theta_{i-\frac{1}{2}} \sum_{\alpha=1}^d R_i^\alpha \left(-\frac{\Delta x}{2}\right)^\alpha \quad \text{and} \quad \tilde{W}_{i,+}^n = W_i^n + \theta_{i+\frac{1}{2}} \sum_{\alpha=1}^d R_i^\alpha \left(\frac{\Delta x}{2}\right)^\alpha,$$

and the new high-order well-balanced source term reads:

$$\tilde{S}_i^n = \left(1 - \frac{\theta_{i-\frac{1}{2}}^n + \theta_{i+\frac{1}{2}}^n}{2}\right) S_i^n + \frac{\theta_{i-\frac{1}{2}}^n + \theta_{i+\frac{1}{2}}^n}{2} \hat{S}_i^n.$$

**Next step:** Provide a suitable choice of the **convex combination parameter  $\theta_{i\pm\frac{1}{2}}$** . We follow the general strategy from [C. Berthon, S. Bulteau, F. Foucher, M. M'Baye and V. M.-D., *SIAM SISC*, 2022].

# Steady solution detector

The convex combination parameter  $\theta_{i+1/2}^n$  must satisfy the following properties:

- vanish when  $(W_i^n, W_{i+1}^n)$  are at equilibrium;
- be an approximation of 1 up to  $\mathcal{O}(\Delta x^{d+1})$  otherwise.

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- vanish when  $(W_i^n, W_{i+1}^n)$  are at equilibrium;
- be an approximation of 1 up to  $\mathcal{O}(\Delta x^{d+1})$  otherwise.

We propose the following expression:

$$\theta_{i+1/2}^n = \frac{\varepsilon_{i+1/2}^n}{\varepsilon_{i+1/2}^n + C_{i+1/2}^n \Delta x^{d+1}},$$

with  $\varepsilon_{i+1/2}^n = \left\| \begin{pmatrix} q_{i+1}^n - q_i^n \\ B(h_{i+1}^n, q_{i+1}^n, Z_{i+1}) - B(h, q_i^n, Z_i) \end{pmatrix} \right\|.$

# Properties of the steady solution detector

$$\theta_{i+\frac{1}{2}}^n = \frac{\varepsilon_{i+\frac{1}{2}}^n}{\varepsilon_{i+\frac{1}{2}}^n + C_{i+\frac{1}{2}}^n \Delta x^{d+1}}, \text{ with } \varepsilon_{i+\frac{1}{2}}^n = \left\| \begin{pmatrix} q_{i+1}^n - q_i^n \\ B(h_{i+1}^n, q_{i+1}^n, Z_{i+1}) - B(h_i^n, q_i^n, Z_i) \end{pmatrix} \right\|$$

(WB) We easily note that  $\varepsilon_{i+\frac{1}{2}}^n$  vanishes (and therefore  $\theta_{i+\frac{1}{2}}^n$  does too) as soon as  $W_i^n$  and  $W_{i+1}^n$  are at equilibrium.

(HO) If  $\varepsilon_{i+\frac{1}{2}}^n \neq 0$ , then

$$\theta_{i+\frac{1}{2}}^n = \frac{1}{1 + \Delta x^{d+1} \frac{C_{i+\frac{1}{2}}^n}{\varepsilon_{i+\frac{1}{2}}^n}} = 1 + \mathcal{O}(\Delta x^{d+1}).$$

$\rightsquigarrow$  The expression of  $\theta_{i\pm\frac{1}{2}}^n$  satisfies the required properties.

**Next step:** perform numerical tests to validate the method.

# Limitations of the method

Of course, the method also has a few limitations.

1. It is dependent on a **parameter**  $C$ , which could be different for each experiment.
2. Although the scheme is high-order accurate and well-balanced, there is an **issue with high-order well-balanced initialization**.

Consider an initial condition  $W_0$ , steady at interface  $x_{i-1/2}$  and unsteady at interface  $x_{i+1/2}$ ; we need the reconstruction  $\tilde{W}_i^0$  to satisfy

$$\tilde{W}_i^0(x_{i-\frac{1}{2}}) = \frac{1}{\Delta x} \int_{x_{i-\frac{1}{2}}}^{x_{i+\frac{1}{2}}} W_0(x) dx \quad \text{and} \quad \tilde{W}_i^0(x_{i+\frac{1}{2}}) = W_0(x_{i+\frac{1}{2}}) + \mathcal{O}(\Delta x^{d+1}).$$

This leads to two conditions in cell  $i$ , for one unknown  $W_i^0$ ...

## An expression of $C_{i+1/2}^n$

To implement the scheme, we need to give an expression of  $C = C_{i+1/2}^n$ .  
We propose  $C_{i+1/2}^0 = 1$ , and, for  $n \geq 1$ :

$$C_{i+\frac{1}{2}}^n = C_\theta \frac{1}{2} \left( \frac{\|W_{i+1}^n - W_{i+1}^{n-1}\|}{\Delta t} + \frac{\|W_i^n - W_i^{n-1}\|}{\Delta t} \right),$$

with  $C_\theta$  a constant parameter.

Note that

$$\theta_{i+\frac{1}{2}}^n = \frac{\varepsilon_{i+\frac{1}{2}}^n}{\varepsilon_{i+\frac{1}{2}}^n + \left( \frac{\Delta x}{C_{i+\frac{1}{2}}^n} \right)^\delta} = \frac{\varepsilon_{i+\frac{1}{2}}^n (C_{i+\frac{1}{2}}^n)^\delta}{\varepsilon_{i+\frac{1}{2}}^n (C_{i+\frac{1}{2}}^n)^\delta + \Delta x^\delta} :$$

we get  $\theta_{i+\frac{1}{2}}^n = 0$  if  $\varepsilon_{i+\frac{1}{2}}^n = 0$  or  $C_{i+\frac{1}{2}}^n = 0$ . Why does this make sense?



## An expression of $C_{i+1/2}^n$ – reasoning

$$\theta_{i+\frac{1}{2}}^n = 0 \text{ if } \varepsilon_{i+\frac{1}{2}}^n = 0 \text{ or } C_{i+\frac{1}{2}}^n = 0$$

$\varepsilon_{i+\frac{1}{2}}^n = 0 \implies$  steady state solution for the equations  
 $\implies \theta_{i+\frac{1}{2}}^n$  must vanish to preserve the steady state solution

$C_{i+\frac{1}{2}}^n = 0 \implies$  vanishing discrete time derivative  
 $\implies$  steady state solution for the high-order scheme  
 $\implies$  not a steady state solution for the equations<sup>1</sup>  
 $\implies \theta_{i+\frac{1}{2}}^n$  must vanish to perturb the solution

---

<sup>1</sup>Otherwise, the high-order scheme would be well-balanced.