Band topology and submanifolds of matrices

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Theme of my talk

Submanifold in the space of matrices

- I will relate:
 - cohomology classes on such submanifolds;
 - topological invariants of quantum systems on lattices such as insulators and semimetals.

What I will talk about is something like a basic idea for future possible works, and it stemmed from discussion with collaborators:

Ken Shiozaki, Masatoshi Sato, Guo Chuan Thiang, ...

Plan of my talk

Introduction

- **②** Submanifolds in the space of matrices
- **③** Change of topological numbers
- Semimetal

Quantum system on lattice

- I would like to consider a certain quantum mechanical system on a lattice Z^d ⊂ R^d.
- The Hilbert space in this system is

$$L^2(\mathbb{Z}^d,\mathbb{C}^n)=iggl\{\psi=(\psi(j))_{j\in\mathbb{Z}^d}igg|\ \sum\|\psi(j)\|^2<\inftyiggr\},$$

where n is the internal freedom at each site.

- The Hamiltonian $H: L^2(\mathbb{Z}^d, \mathbb{C}^n) \to L^2(\mathbb{Z}^d, \mathbb{C}^n)$ is a self-adjoint operator which is
 - commuting with the translation operator on the lattice,
 - describing a "short range interaction".

Fourier transformed Hamiltonian

• Under the assumptions on the Hamiltonian *H*, its information is completely encoded into the Fourier transformed Hamiltonian

$$\hat{H}: T^d \to \mathcal{H}(\mathbb{C}^n).$$

This is a continuous map from the *d*-dimensional torus (BZ torus) to the space of $n \times n$ Hermitian matrices

 $\mathcal{H}(\mathbb{C}^n) = \{ H \in M(n,\mathbb{C}) | \ H^{\dagger} = H \} \cong \mathbb{R}^{n^2}.$

- Let us say that the quantum system is
 - insulator if $\det \hat{H}(k) \neq 0$ for all $k \in T^d$,
 - semimetal if $\det \hat{H}(k) = 0$ at some $k \in T^d$.

(The definition may be physically insufficient.)

Classification of topological phase by homotopy

- Now, the classification of the topological phase described by *Ĥ* can be done by the classification of homotopy class of *Ĥ* within a fixed type of phases.
- For example, if \hat{H} is an insulator $(\det \hat{H} \neq 0)$, then we consider the homotopy within insulators. In other words, we consider the homotopy classes of

$$T^d o \mathcal{H}_0(\mathbb{C}^n) = \{ H \in \mathcal{H}(\mathbb{C}^n) | \det H \neq 0 \}.$$

 A more equivalence relation (such as "stability" under an addition of uninteresting phases) can be introduced, but I will not consider it today for simplicity.

Topological invariant: the case of insulator

- Then, for the classification of the homotopy class of \hat{H} , it is useful to consider a topological invariant of \hat{H} , a quantity invariant under the homotopy deformation.
- For example, in the case of an insulator

$$\hat{H}: T^d o \mathcal{H}_0(\mathbb{C}^n) = \{ H \in \mathcal{H}(\mathbb{C}^n) | \ \det H
eq 0 \},$$

one has the Chern classes of the Bloch vector bundle

$$E = igcup_{k\in T^d} igcup_{\lambda<0} \operatorname{Ker}(\hat{H}(k)-\lambda).$$

• In 2D case, the integration of the Chern class $c_1(E)$ is a topological invariant of the insulator.

Topological invariant: the case of insulator

• In the example of an insulator, the Chern classes of Bloch bundles are closely related to the subspace

 $\mathcal{H}_0(\mathbb{C}^n) = \{ H \in \mathcal{H}(\mathbb{C}^n) | \det H \neq 0 \}$

in the space of Hermitian matrices $\mathcal{H}(\mathbb{C}^n)$.

- Actually, $\mathcal{H}_0(\mathbb{C}^n) \subset \mathcal{H}(\mathbb{C}^n)$ is an open submanifold homotopy equivalent to the Grassmannian of \mathbb{C}^n .
- To consider semimetals, we allow singularities for the Hamiltonian $\hat{H}: T^d \to \mathcal{H}(\mathbb{C}^n)$, namely, points $k \in T^d$ at which $\det \hat{H}(k) = 0$.
- This leads us to consider subspaces in $\mathcal{H}(\mathbb{C}^n)$ other than $\mathcal{H}_0(\mathbb{C}^n)$.

The things I would like to talk today:

- **1** Some subspaces of matrices give rise to submanifolds.
- Some cohomology classes on such submanifolds provide invariants of insulators and semimetals.
- An application is a formula of the change of topological numbers under a deformation allowing singularities.

The things I would like to talk today:

- **1** Some subspaces of matrices give rise to submanifolds.
- Some cohomology classes on such submanifolds provide invariants of insulators and semimetals.
- An application is a formula of the change of topological numbers under a deformation allowing singularities.
 - Originally, I anticipated that a use of the submanifolds of matrices can be applied to a detection of new topological phases. But, the things I tell you today are still on the way to the goal.

Introduction

- **②** Submanifolds in the space of matrices
- **O Change of topological numbers**
- Semimetal

The space of complex square matrices

• To warm up, I would like to consider submanifolds in the space of complex matrices $\mathcal{M}(\mathbb{C}^n) = M(n, \mathbb{C})$, rather than those in the space of Hermitian matrices $\mathcal{H}(\mathbb{C}^n)$.

The space of complex square matrices

- To warm up, I would like to consider submanifolds in the space of complex matrices $\mathcal{M}(\mathbb{C}^n) = M(n, \mathbb{C})$, rather than those in the space of Hermitian matrices $\mathcal{H}(\mathbb{C}^n)$.
- If we consider a quantum system of class AIII, then the chiral symmetry arrows us to assume that a Fourier-transformed Hamiltonian is of the form

$$\hat{H}(k) = \left(egin{array}{cc} 0 & A(k) \ A(k)^{\dagger} & 0 \end{array}
ight).$$

Hence the Hamiltonian \hat{H} amounts to

$$A: T^d o \mathcal{M}(\mathbb{C}^n) = M(n, \mathbb{C}).$$

• Because of this fact, it is still meaningful to study $\mathcal{M}(\mathbb{C}^n)$ in the context of quantum systems.

Subspaces in $\mathcal{M}(\mathbb{C}^n) = M(n,\mathbb{C}) \cong \mathbb{C}^{n^2}$

Definition

For an integer k, we put

$$\mathcal{M}_{\leq k}(\mathbb{C}^n) := \{A \in \mathcal{M}(\mathbb{C}^n) | \dim \operatorname{Ker} A \leq k\}, \ \mathcal{M}_k(\mathbb{C}^n) := \{A \in \mathcal{M}(\mathbb{C}^n) | \dim \operatorname{Ker} A = k\}.$$

Here A is regarded as a linear map $A: \mathbb{C}^n \to \mathbb{C}^n$ by the left multiplication.

$$\mathcal{M}_{\leq k}(\mathbb{C}^n) = \mathcal{M}_{\leq k-1}(\mathbb{C}^n) \cup \mathcal{M}_k(\mathbb{C}^n)$$
$$\cdots \subset \mathcal{M}_{\leq k-1}(\mathbb{C}^n) \subset \mathcal{M}_{\leq k}(\mathbb{C}^n) \subset \cdots$$

• In general, $\mathcal{M}_{\leq k}(\mathbb{C}^n) \subset \mathcal{M}(\mathbb{C}^n)$ is an open complex submanifold, and $\dim_{\mathbb{C}} \mathcal{M}_{< k} = \dim_{\mathbb{C}} \mathcal{M} = n^2$.

Example: n = 1

$$\mathcal{M}_{\leq k}(\mathbb{C}^n) = \{A \in \mathcal{M}(\mathbb{C}^n) | \dim \operatorname{Ker} A \leq k\} \subset^{\operatorname{open}} \mathcal{M}(\mathbb{C}^n), \ \mathcal{M}_k(\mathbb{C}^n) = \{A \in \mathcal{M}(\mathbb{C}^n) | \dim \operatorname{Ker} A = k\}.$$

$$\begin{split} \mathcal{M}_{\leq 1}(\mathbb{C}) &= M(1,\mathbb{C}) = \mathbb{C} \\ & \bigcup \qquad \qquad \mathcal{M}_{\leq 0}(\mathbb{C}) &= GL(1,\mathbb{C}) = \mathbb{C}^{\times} \end{split}$$

• In general, we have obvious identifications

$$egin{aligned} \mathcal{M}_{\leq n}(\mathbb{C}^n) &= \mathcal{M}(\mathbb{C}^n) \cong \mathbb{C}^{n^2}, & \mathcal{M}_n(\mathbb{C}^n) = \{0\}, \ \mathcal{M}_{\leq n-1}(\mathbb{C}^n) &\cong \mathbb{C}^{n^2} ackslash \{0\}, \ \mathcal{M}_{\leq 0}(\mathbb{C}^n) &= \mathcal{M}_0(\mathbb{C}^n) = GL(n,\mathbb{C}). \end{aligned}$$

Example: n = 2

$$\mathcal{M}_{\leq k}(\mathbb{C}^n) = \{A \in \mathcal{M}(\mathbb{C}^n) | \dim \operatorname{Ker} A \leq k\} \stackrel{\mathsf{open}}{\subset} \mathcal{M}(\mathbb{C}^n), \ \mathcal{M}_k(\mathbb{C}^n) = \{A \in \mathcal{M}(\mathbb{C}^n) | \dim \operatorname{Ker} A = k\}.$$

$$egin{aligned} \mathcal{M}_{\leq 2}(\mathbb{C}^2)&=M(2,\mathbb{C})\cong\mathbb{C}^4\ &igcup\ &egcup\ &igcup\ &igcup\ &$$

Subspace $\mathcal{M}_k(\mathbb{C}^n) \subset \mathcal{M}_{\leq k}(\mathbb{C}^n)$

$$\mathcal{M}_{\leq k}(\mathbb{C}^n) = \{A \in \mathcal{M}(\mathbb{C}^n) | \dim \operatorname{Ker} A \leq k\} \stackrel{\mathsf{open}}{\subset} \mathcal{M}(\mathbb{C}^n), \\ \mathcal{M}_k(\mathbb{C}^n) = \{A \in \mathcal{M}(\mathbb{C}^n) | \dim \operatorname{Ker} A = k\}.$$

Fact

 $\mathcal{M}_k(\mathbb{C}^n)\subset\mathcal{M}_{\leq k}(\mathbb{C}^n)$ is a closed complex submanifold of

$$\dim_{\mathbb{C}}\mathcal{M}_k(\mathbb{C}^n)=n^2-k^2.$$

$$egin{aligned} \mathcal{M}_{\leq 1}(\mathbb{C}) &= M(1,\mathbb{C}) = \mathbb{C} \ & igcup \ & igcup \ & \mathcal{M}_{1}(\mathbb{C}) = \{0\} \ & \mathcal{M}_{\leq 0}(\mathbb{C}) &= GL(1,\mathbb{C}) = \mathbb{C}^{ imes} \end{aligned}$$

Example: n = 2

$$\mathcal{M}_{\leq k}(\mathbb{C}^n) = \{A \in \mathcal{M}(\mathbb{C}^n) | \dim \operatorname{Ker} A \leq k\} \stackrel{\mathsf{op}}{\subset} \mathcal{M}(\mathbb{C}^n), \ \mathcal{M}_k(\mathbb{C}^n) = \{A \in \mathcal{M}(\mathbb{C}^n) | \dim \operatorname{Ker} A = k\} \underset{\mathsf{cl}}{\subset} \mathcal{M}_{\leq k}(\mathbb{C}^n).$$

$$egin{aligned} \mathcal{M}_{\leq 2}(\mathbb{C}^2) &= M(2,\mathbb{C}) \cong \mathbb{C}^4 \ & igcup & \mathcal{M}_2(\mathbb{C}^2) = \{0\} \ \mathcal{M}_{\leq 1}(\mathbb{C}^2) &\cong \mathbb{C}^4 igcarrow \{0\} \ & igcup & \mathcal{M}_1(\mathbb{C}^2) = \{A
eq 0, \det A = 0\} \ \mathcal{M}_{\leq 0}(\mathbb{C}^2) &= GL(2,\mathbb{C}) \end{aligned}$$

Cohomology

$$\begin{array}{ll} \mathcal{M}_{\leq k}(\mathbb{C}^n) \ = \mathcal{M}_{\leq k-1}(\mathbb{C}^n) \cup \mathcal{M}_k(\mathbb{C}^n) \ \supset \ \mathcal{M}_k(\mathbb{C}^n) \\ \bigcup & \text{closed} \\ \mathcal{M}_{\leq k-1}(\mathbb{C}^n) \text{ open} \end{array}$$

- To apply the submanifolds of matrices, we need to know their cohomology groups.
- A basic tool is the long exact sequence associated to the pair $(\mathcal{M}_{\leq k}(\mathbb{C}^n), \mathcal{M}_{\leq k-1}(\mathbb{C}^n))$:
- $\cdots o H^m(\mathcal{M}_{\leq k}, \mathcal{M}_{\leq k-1}) o H^m(\mathcal{M}_{\leq k}) o H^m(\mathcal{M}_{\leq k-1}) o$ $\cdots o H^{m+1}(\mathcal{M}_{\leq k}, \mathcal{M}_{\leq k-1}) o \cdots$

Cohomology

$$\begin{array}{ll} \mathcal{M}_{\leq k}(\mathbb{C}^n) \ = \mathcal{M}_{\leq k-1}(\mathbb{C}^n) \cup \mathcal{M}_k(\mathbb{C}^n) \ \supset \ \mathcal{M}_k(\mathbb{C}^n) \\ \bigcup & \text{closed} \\ \mathcal{M}_{\leq k-1}(\mathbb{C}^n) \text{ open} \end{array}$$

• Another basic tool is an isomorphism ("Alexander-Poincaré duality")

 $H^m(\mathcal{M}_{\leq k}(\mathbb{C}^n),\mathcal{M}_{\leq k-1}(\mathbb{C}^n)) \stackrel{lpha}{\cong} H^{m-2k^2}(\mathcal{M}_k(\mathbb{C}^n)).$

• n = 1 and k = 1:

$$egin{aligned} &H^m(\mathcal{M}_{\leq 1},\mathcal{M}_{\leq 0})=H^m(\mathbb{C},\mathbb{C}^{ imes})\cong H^m(D^2,\partial D^2),\ &H^{m-2}(\mathcal{M}_1)=H^{m-2}(\{0\}). \end{aligned}$$

Remark

- Next, I will apply the submanifolds and their cohomology classes to quantum systems of class AIII.
- Before that, I remark that what I presented here is a finite-dimensional analogue of Koschorke's work:

U. Koschorke, Infinite dimensional *K*-theory and characteristic classes of Fredholm bundle maps. 1970 Global Analysis (Proc. Sympos. Pure Math., Vol. XV, Berkeley, Calif., 1968) pp. 95–133 Amer. Math. Soc., Providence, R.I.

• In this paper, characteristic classes for *K*-theory are constructed by using submanifolds in the space of Fredholm operators.

Remark

Replacing Cⁿ by an infinite-dimensional Hilbert space
 ℋ, we can make sense of the following generalizations of
 ℳ_{≤k}(Cⁿ) and ℳ_k(Cⁿ):

$$\mathcal{F}_{\leq k}(\mathcal{H}) = \left\{egin{array}{l} A: \mathcal{H}
ightarrow \mathcal{H} extsf{Fredholm operator} \ \dim \mathrm{Ker} A = \dim \mathrm{Coker} A \leq k \end{array}
ight\}, \ \mathcal{F}_k(\mathcal{H}) = \left\{egin{array}{l} A: \mathcal{H}
ightarrow \mathcal{H} extsf{Fredholm operator} \ \dim \mathrm{Ker} A = \dim \mathrm{Coker} A = k \end{array}
ight\}.$$

- It turns out that $\mathcal{F}_k(\mathcal{H}) \subset \mathcal{F}_{\leq k}(\mathcal{H})$ is a complex submanifold of codimension k^2 , although $\mathcal{F}_k(\mathcal{H})$ and $\mathcal{F}_{\leq k}(\mathcal{H})$ are infinite-dimensional.
- There is also an isomorphism

$$H^m(\mathcal{F}_{\leq k}, \mathcal{F}_{\leq k-1}) \cong H^{m-2k^2}(\mathcal{F}_k).$$

Introduction

2 Submanifolds in the space of matrices

O Change of topological numbers

Semimetal

Change of topological numbers

- Recall that a certain quantum system on a *d*-dimensional lattice is described by a map $H: T^d \to \mathcal{H}(\mathbb{C}^n)$.
- An insulator is characterized by the gap condition; $\det H(k) \neq 0$ for all $k \in T^d$.
- A quantum system is class AIII, if there is a linear symmetry Γ (called a chiral symmetry) such that

$$\Gamma^2 = 1,$$
 $\Gamma H(k) = -H(k)\Gamma.$

• Without loss of generality, we can assume

$$\Gamma = \left(egin{array}{cc} 1_{\mathbb{C}^n} & 0 \ 0 & -1_{\mathbb{C}^n} \end{array}
ight), \ \ H(k) = \left(egin{array}{cc} 0 & A(k)^\dagger \ A(k) & 0 \end{array}
ight).$$

Class AllI insulator

• Then a *d*-dim insulator of class AIII is described by

 $A:T^d
ightarrow\mathcal{M}_{\leq 0}(\mathbb{C}^n)=\mathcal{M}_0(\mathbb{C}^n)=GL(n,\mathbb{C}).$

 ${\ }$ It is well-known that $GL(n,\mathbb{C})\simeq U(n),$ and

$$H^{ullet}(U(n);\mathbb{Z})=igwedge(
u_1,\dots,
u_{2n-1}),$$

where $u_{2j-1} \in H^{2j-1}(U(n);\mathbb{Z}).$

• In the case that d = 2n - 1, we can define the *d*-dimensional winding number by

$$\operatorname{top}(A) = \int_{T^{2n-1}} A^*
u_{2n-1} \ \in \mathbb{Z}.$$

Change of topological phases

• Now, let us consider a one-parameter family of (2n-1)-dimensional class AIII quantum systems

$$ilde{A}:T^{2n-1} imes [0,1] o \mathcal{M}(\mathbb{C}^n)=\mathbb{C}^{n^2}.$$

Assumption

- **()** \tilde{A} is a smooth map.
- $\widehat{A}(k,t) \in \mathcal{M}_{\leq 1}(\mathbb{C}^n) \subset \mathcal{M}(\mathbb{C}^n) \text{ for all } k,t.$
- $\tilde{A}(k,i) \in \mathcal{M}_0(\mathbb{C}^n)$ for all k and i = 0, 1.
- a condition which ensures that

$$\Sigma := \{(k,t) \in T^{2n-1} imes [0,1] | \det ilde{A}(k,t) = 0\}$$

is a (2n-2)-dimensional submanifold of $T^{2n-1} \times [0,1]$.

Change of topological numbers

$$egin{aligned} &(T^{2n-1} imes [0,1],T^{2n-1} imes \{0,1\})\stackrel{ ilde{A}}{\longrightarrow} (\mathcal{M}_{\leq 1}(\mathbb{C}^n),\mathcal{M}_{\leq 0}(\mathbb{C}^n))\ &\Sigma=\{(k,t)\in T^{2n-1} imes [0,1]| \ \det ilde{A}(k,t)=0\}\stackrel{ ilde{A}ert_{\Sigma}}{\longrightarrow} \mathcal{M}_1(\mathbb{C}^n) \end{aligned}$$

Theorem

Under the assumption, we have

$$ext{top}(ilde{A}|_{T^{2n-1} imes \{0\}}) - ext{top}(ilde{A}|_{T^{2n-1} imes \{1\}}) = \int_{\Sigma} ilde{A}|_{\Sigma}^* \gamma_{2n-2},$$

where $\gamma_{2n-2} \in H^{2n-2}(\mathcal{M}_1(\mathbb{C}^n))$ is the class induced from $\nu_{2n-1} \in H^{2n-1}(\mathcal{M}_{\leq 0}(\mathbb{C}^n))$ through

$$H^{2n-1}(\mathcal{M}_{\leq 0}) \stackrel{\delta}{
ightarrow} H^{2n}(\mathcal{M}_{\leq 1}, \mathcal{M}_{\leq 0}) \stackrel{lpha}{\cong} H^{2n-2}(\mathcal{M}_1).$$

Example: d = 1 (n = 1)

$$egin{aligned} &(T^{2n-1} imes [0,1],T^{2n-1} imes \{0,1\})\stackrel{ ilde{A}}{\longrightarrow} (\mathcal{M}_{\leq 1}(\mathbb{C}^n),\mathcal{M}_{\leq 0}(\mathbb{C}^n))\ &\Sigma=\{(k,t)\in T^{2n-1} imes [0,1]|\;\det ilde{A}(k,t)=0\}\ & ext{top}(ilde{A}|_{T^{2n-1} imes \{0\}})- ext{top}(ilde{A}|_{T^{2n-1} imes \{1\}})=\int_{\Sigma} ilde{A}|_{\Sigma}^{*}\gamma_{2n-2} \end{aligned}$$

• In the case of n = 1, let us consider

$$egin{array}{lll} ilde{A}: \ \mathbb{R}/2\pi\mathbb{Z} imes [0,1] \longrightarrow \mathcal{M}_{\leq 1}(\mathbb{C}) = \mathbb{C} \ ilde{A}(k,t) = (1-t) + te^{ik} \end{array}$$

• Then the change of the topological numbers

$$0 - 1 = -1$$

agrees with the number of points in $\Sigma = \{(\pi, 1/2)\}$ counted with weight $\tilde{A}^* \gamma_0 = -1$.

Remark

- I just explained the case of class AIII, but the formula can be generalized:
 - we can consider the other Altland-Zirnbauer classes (class A, AI, BDI, D,)
 - We can allow crystalline symmetry of the lattice
- In the first generalization, we need to consider a suitable equivariant cohomology, for real classes.
- In the second generalization, we need to consider a fiber bundle of submanifolds of matrices in addition, for nonsymmorphic crystalline symmetry.

Introduction

- **②** Submanifolds in the space of matrices
- **Or Change of topological numbers**

Semimetal

• I would like to relate submanifolds of matrices with the "semimetal class" of V. Mathai and G. C. Thiang:

V. Mathai and G. C. Thiang, Global topology of Weyl semimetals and Fermi arcs. Journal of Physics A: Mathematical and Theoretical 50 (11), 11LT01

V. Mathai and G. C. Thiang, Differential topology of semimetals. Communications in Mathematical Physics 355 (2), 561-602

Submanifolds of traceless Hermitian matrices

• For the semimetal setup, let us consider the subspaces of $n \times n$ traceless Hermitian matrices

$$\mathcal{H}^{\mathbf{0}}(\mathbb{C}^n) = \{H \in \mathcal{M}(\mathbb{C}^n) | \ H^{\dagger} = H, \operatorname{tr} H = \mathbf{0}\} \cong \mathbb{R}^{n^2 - 1},$$

 $\mathcal{H}^{\mathbf{0}}_{\leq k}(\mathbb{C}^n) = \{H \in \mathcal{H}^{\mathbf{0}}(\mathbb{C}^n) | \ \operatorname{dim}\operatorname{Ker} H \leq k\},$
 $\mathcal{H}^{\mathbf{0}}_k(\mathbb{C}^n) = \{H \in \mathcal{H}^{\mathbf{0}}(\mathbb{C}^n) | \ \operatorname{dim}\operatorname{Ker} H = k\}.$

• The traceless condition leads to

$$\mathcal{H}^0_{n-1}(\mathbb{C}^n)=\emptyset, \hspace{0.5cm} \mathcal{H}^0_{\leq n-2}(\mathbb{C}^n)=\mathcal{H}^0_{\leq n-1}(\mathbb{C}^n).$$

• Except for this point, we have the same results as before.

Submanifolds of traceless Hermitian matrices

$$egin{aligned} \mathcal{H}^0_{\leq k}(\mathbb{C}^n) &= \{ H \in \mathcal{H}^0(\mathbb{C}^n) | \ \mathrm{dim}\mathrm{Ker} H \leq k \}, \ \mathcal{H}^0_k(\mathbb{C}^n) &= \{ H \in \mathcal{H}^0(\mathbb{C}^n) | \ \mathrm{dim}\mathrm{Ker} H = k \}. \end{aligned}$$

- In general, $\mathcal{H}^0_{\leq k}(\mathbb{C}^n) \subset \mathcal{H}^0(\mathbb{C}^n)$ is an open submanifold of dimension $\dim_{\mathbb{R}} \mathcal{H}^0_{\leq k}(\mathbb{C}^n) = \dim_{\mathbb{R}} \mathcal{H}^0(\mathbb{C}^n) = n^2 1$.
- $\mathcal{H}_k^0(\mathbb{C}^n) \subset \mathcal{H}_{\leq k}^0(\mathbb{C}^n)$ is a closed submanifold of dimension $n^2 k^2 1$ for $k \leq n 2$, with orientable normal bundle.
- In the case of n = 2:

$$egin{aligned} \mathcal{H}^0_{\leq 2}(\mathbb{C}^2)&\cong \mathbb{R}^3\ &igcup\ &egcup\ &igcup\ &egcup\ &egu\ &igcup\ &igcup\ &$$

3d "Weyl" semimetal

$$\begin{split} & \mathcal{H}^0_{\leq 2}(\mathbb{C}^2) \cong \mathbb{R}^3 \\ & \bigcup \\ & \mathcal{H}^0_{\leq 0}(\mathbb{C}^2) = \mathcal{H}^0_{\leq 1}(\mathbb{C}^2) \cong \mathbb{R}^3 \backslash \{0\} \\ \end{split}$$

• Suppose that we are given a continuous map

$$H: T^3
ightarrow \mathcal{H}^0_{\leq 2}(\mathbb{C}^2) = \mathcal{H}^0(\mathbb{C}^2) \cong \mathbb{R}^3$$

such that the set of "Weyl points"

 $W = \{k \in T^3 | \det H(k) = 0\}$

consists of a finite number of points.

3d "Weyl" semimetal

- $egin{aligned} W &= \{k \in T^3 | \ \det H(k) = 0\} \ W & \mathcal{H}^0_2(\mathbb{C}^2) = \{0\} \ \cap & \cap \ T^3 & \stackrel{H}{ o} & \mathcal{H}^0_{\leq 2}(\mathbb{C}^2) \cong \mathbb{R}^3 \ \cup & \cup \ T^3 igvee W & \mathcal{H}^0_{\leq 0}(\mathbb{C}^2) = \mathcal{H}^0_{\leq 1}(\mathbb{C}^2) \cong \mathbb{R}^3 igvee 0\} \simeq S^2 \end{aligned}$
- We can then have a map

$$H: T^3 \setminus W \to \mathcal{H}^0_{\leq 0}(\mathbb{C}^2).$$

• The pullback of a generator in $H^2(\mathcal{H}^0_{\leq 0}(\mathbb{C}^2)) \cong \mathbb{Z}$ is essentially the semimetal class of Mathai and Thiang:

The semimetal class of $H \in H^2(T^3 \setminus W)$.

3d "Weyl" semimetal

• The "Mayer-Vietoris sequence" of Mathai-Thiang is equivalent to

$$\overbrace{H^2(T^3)}^{\text{insulator}} \to \overbrace{H^2(T^3 \backslash W)}^{\text{semimetal}} \stackrel{\delta}{\to} H^3(T^3, T^3 \backslash W) \stackrel{\alpha}{\cong} \overbrace{H^0(W)}^{\text{charge}}.$$

• There is the universal counterpart

$$H^2(\mathcal{H}^0_{\leq 2}) o H^2(\mathcal{H}^0_{\leq 0}) \stackrel{\delta}{ o} H^3(\mathcal{H}^0_{\leq 2},\mathcal{H}^0_{\leq 0}) \stackrel{lpha}{\cong} H^0(\mathcal{H}^0_2).$$

• A use of the submanifolds

$$\cdots \subset \mathcal{H}^0_{\leq k-1}(\mathbb{C}^n) \subset \mathcal{H}^0_{\leq k}(\mathbb{C}^n) \subset \cdots$$

would generalize the semimetal class.